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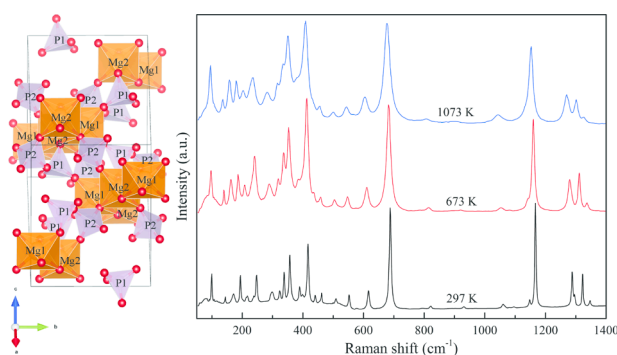
## Temperature-dependent Raman spectra and thermal expansion of $\text{MgP}_2\text{O}_6$

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### HIGHLIGHTS

- Raman spectra and X-ray diffraction patterns of  $\text{MgP}_2\text{O}_6$  were collected up to 1073 K.
- No temperature-induced phase transition occurs for  $\text{MgP}_2\text{O}_6$ .
- The effect of temperature on the Raman active vibrations of  $\text{MgP}_2\text{O}_6$  was analyzed.
- The thermal expansion coefficients and isobaric mode Grüneisen parameters of  $\text{MgP}_2\text{O}_6$  were determined.

### GRAPHICAL ABSTRACT



### ARTICLE INFO

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$\text{MgP}_2\text{O}_6$   
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### ABSTRACT

The Raman spectra and thermal expansion of  $\text{MgP}_2\text{O}_6$  metaphosphate were investigated at various temperatures up to 1073 K at ambient pressure. No temperature-induced phase transition was observed in this study. The effect of temperature on the Raman active vibrations and unit cell parameters was determined. All the observed Raman active bands of  $\text{MgP}_2\text{O}_6$  showed linear temperature dependences in the range of  $-2.61 \times 10^{-2} \sim -0.49 \times 10^{-2} \text{ cm}^{-1} \text{ K}^{-1}$ . The thermal expansion coefficient of  $\text{MgP}_2\text{O}_6$  was estimated to be  $3.21(2) \times 10^{-5} \text{ K}^{-1}$ . An axial anisotropic thermal expansion exists and the *c*-axis shows the smallest thermal expansion. The isobaric mode Grüneisen parameters of  $\text{MgP}_2\text{O}_6$  were calculated. The obtained results were compared with other compounds in the  $\text{MgO-P}_2\text{O}_5$  system.

### 1. Introduction

The metaphosphates with the chemical formula of  $\text{AP}_2\text{O}_6$  (A = divalent cation) contain a number of compounds. Due to their potential technical and material applications, such as luminescence, dielectric,

biomaterial, and catalysis, metaphosphates have been intensively investigated [1–10]. Mainly there are three kinds of crystal structures for  $\text{AP}_2\text{O}_6$  metaphosphates with different divalent cations, including monoclinic  $\text{CuP}_2\text{O}_6$ -type (space group:  $C2/c$ , e.g., A = Mg, Cu, Ni, Co, Fe, Mn) and  $\text{CaP}_2\text{O}_6$ -type (space group:  $P2_1/c$ , e.g., A = Ca, Sr, Ba, Pb), and

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orthorhombic CdP<sub>2</sub>O<sub>6</sub>-type (space group: *Pbca*, e.g., A = Cd, Hg) [11]. Additionally, a high-temperature form of CdP<sub>2</sub>O<sub>6</sub> also belongs to the orthorhombic structure (space group: *P2<sub>1</sub>2<sub>1</sub>2<sub>1</sub>*) [11].

Magnesium metaphosphate, MgP<sub>2</sub>O<sub>6</sub>, is one of the four compounds in MgO-P<sub>2</sub>O<sub>5</sub> binary system [12–13]. The crystal structure of monoclinic MgP<sub>2</sub>O<sub>6</sub> was determined in previous studies [14–15], as shown in Fig. 1. In MgP<sub>2</sub>O<sub>6</sub> two kinds of Mg with six-fold coordination exist and slightly disordered MgO<sub>6</sub> octahedra in the interlayer space share edges, and two kinds of P with four-fold coordination form corner-sharing PO<sub>4</sub> tetrahedra, which are linked together in a ring to form P<sub>4</sub>O<sub>12</sub><sup>4-</sup> anions by P-O-P bridges [15]. The three-dimensional framework of MgP<sub>2</sub>O<sub>6</sub> is constructed by P<sub>4</sub>O<sub>12</sub><sup>4-</sup> anionic groups connected via MgO<sub>6</sub> octahedra.

In previous studies [2,16–21], the melting points of MgP<sub>2</sub>O<sub>6</sub> were reported around 1438 K by differential thermal analysis and quenching experiments, and no temperature-induced phase transition was observed. However, there is little information about the vibrational and thermal expansion properties of MgP<sub>2</sub>O<sub>6</sub> at various temperatures.

In MgO-P<sub>2</sub>O<sub>5</sub> binary system, there are four compounds including Mg<sub>3</sub>(PO<sub>4</sub>)<sub>2</sub>, Mg<sub>2</sub>P<sub>2</sub>O<sub>7</sub>, MgP<sub>2</sub>O<sub>6</sub> and MgP<sub>4</sub>O<sub>11</sub> [12–13]. In our previous studies, the thermal expansions and temperature-dependent Raman spectra of Mg<sub>3</sub>(PO<sub>4</sub>)<sub>2</sub> and Mg<sub>2</sub>P<sub>2</sub>O<sub>7</sub> have been investigated [22–23]. In this study, by using Raman spectroscopic and in-situ X-ray diffraction (XRD) measurements, the Raman spectra and thermal expansion of MgP<sub>2</sub>O<sub>6</sub> were investigated up to 1073 K at ambient pressure. The effects of temperature on Raman active modes and the lattice parameters of MgP<sub>2</sub>O<sub>6</sub> were quantitatively analyzed. Additionally, the thermal behavior of MgP<sub>2</sub>O<sub>6</sub> was compared with other compounds in the MgO-P<sub>2</sub>O<sub>5</sub> system.

## 2. Experimental

A pure MgP<sub>2</sub>O<sub>6</sub> sample was synthesized by a conventional solid-state chemical reaction at 1273 K for 72 h using proportional high pure MgO and NH<sub>4</sub>H<sub>2</sub>PO<sub>4</sub> reagents (99.99% in mass fraction) purchased from Shanghai Aladdin Biochemical Technology Co., Ltd, China. The synthetic MgP<sub>2</sub>O<sub>6</sub> was characterized by powder X-ray diffraction and confirmed as a single phase.

Raman spectra of MgP<sub>2</sub>O<sub>6</sub> at temperatures up to 1073 K and ambient pressure were collected by a Horiba LabRam HR Evolution Raman spectrometer equipped with an 1800 gr/mm grating. The procedure was same as our previous studies [22–24]. A pulsed YAG: Nd<sup>3+</sup> laser beam

with 532-nm wavelength and 100-mW power was adopted as an excitation source. Through a 100× objective lens, the spot of Raman scattering light on the sample surface was about 1 μm<sup>2</sup>. A thin small piece of MgP<sub>2</sub>O<sub>6</sub> was put on a sapphire window which was placed into an alumina chamber in a Linkam TS 1500 high temperature stage for heating. An S-type thermocouple was used to measure the temperatures which were completely automatically and programmatically controlled to change and maintain any desired temperature. For thermal equilibrium in measurements, each Raman spectrum was recorded after keeping at the desired temperature for at least 5 min. The Raman shift was calibrated by plasma and neon emission lines with an accuracy of ±1 cm<sup>-1</sup>. The accumulation time for each spectrum was 30 s, and the obtained spectrum was the average of three measurements. The Raman spectra were analyzed using the PeakFit program (SPSS Inc., Chicago) to obtain a reasonable approximation.

High-temperature in-situ powder X-ray diffraction experiments were carried out from room temperature to 1073 K at the BL14B1 beam line of Shanghai Synchrotron Radiation Facility (SSRF) [25]. The method was same as our previous studies [22–23]. The synthetic MgP<sub>2</sub>O<sub>6</sub> powder was loaded into a fused quartz capillary (inner diameter of 0.5-mm) which was centered at a custom-made furnace for heating and the temperature was measured by a K-type thermocouple. The wavelength of the X-ray beam (beam size = 180 × 200 μm<sup>2</sup>) with an energy of 18 keV was calibrated by the LaB<sub>6</sub> standard from NIST (660b). In-situ XRD patterns were collected using a Mythen 1 K linear detector in Debye-Scherrer mode [26]. For thermal equilibration, each temperature was kept for about 5 min before measuring. The lattice parameters of MgP<sub>2</sub>O<sub>6</sub> at different temperatures were obtained by using the full profile-fitting technique in EXPGUI/GSAS [27–28].

## 3. Results and discussion

As mentioned above, the crystal structure of monoclinic MgP<sub>2</sub>O<sub>6</sub> is in the space group of *C2/c* with *Z* = 4. According to the factor group analysis [29], the Raman active modes of MgP<sub>2</sub>O<sub>6</sub> are as following:

$$\Gamma = 25A_g + 26B_g.$$

Therefore, totally 51 Raman active modes MgP<sub>2</sub>O<sub>6</sub> are predicted. However, due of the overlapping and/or low intensity of some Raman active modes, the number of observed Raman peaks is less than the expected.

As shown in Fig. 2, 29 bands can be distinguished in the Raman spectrum collected at ambient conditions, which is consistent with the previous study [30]. According to previous studies [10,30], the Raman active modes can be ascribed to the vibrations of the P<sub>4</sub>O<sub>12</sub><sup>4-</sup> anionic groups involving P-O-P and = PO<sub>2</sub>. As listed in Table 1, the observed Raman vibrations of MgP<sub>2</sub>O<sub>6</sub> were assigned based on previous reports [10,30]. It is noted that the vibrations of MgP<sub>2</sub>O<sub>6</sub> are more complicated than those of Mg<sub>2</sub>P<sub>2</sub>O<sub>7</sub> involving P<sub>2</sub>O<sub>7</sub><sup>4-</sup> anions and Mg<sub>3</sub>(PO<sub>4</sub>)<sub>2</sub> involving PO<sub>4</sub><sup>3-</sup> anions. It is known that the crystal structure is an important factor for the observed Raman bands of a material, even if the involving elements are the same but in different arrangement. The compounds in MgO-P<sub>2</sub>O<sub>5</sub> system contain basic Mg-O polyhedra and PO<sub>4</sub> tetrahedra with different arrangements. MgP<sub>2</sub>O<sub>6</sub> is monoclinic with space group of *C2/c* where corner-sharing PO<sub>4</sub> tetrahedra form P<sub>4</sub>O<sub>12</sub><sup>4-</sup> anions by P-O-P bridges and connect via MgO<sub>6</sub> octahedra [15]. α- and β-Mg<sub>2</sub>P<sub>2</sub>O<sub>7</sub> are monoclinic with different space groups of *P2<sub>1</sub>/c* and *C2/m* [31–32]. In α-Mg<sub>2</sub>P<sub>2</sub>O<sub>7</sub> the [P<sub>2</sub>O<sub>7</sub>]<sup>4-</sup> groups lie between edge-sharing MgO<sub>5</sub> and MgO<sub>6</sub> polyhedral sheets parallel to the *ab* plane [31]. In β-Mg<sub>2</sub>P<sub>2</sub>O<sub>7</sub> the [P<sub>2</sub>O<sub>7</sub>]<sup>4-</sup> groups linearly lie between the MgO<sub>6</sub> octahedral sheets [32]. There is a difference for the [P<sub>2</sub>O<sub>7</sub>]<sup>4-</sup> anion in α- and β-Mg<sub>2</sub>P<sub>2</sub>O<sub>7</sub> since the P-O-P bond angle is 144° in α-Mg<sub>2</sub>P<sub>2</sub>O<sub>7</sub> [31] but 180° in β-Mg<sub>2</sub>P<sub>2</sub>O<sub>7</sub> [32]. Mg<sub>3</sub>(PO<sub>4</sub>)<sub>2</sub>-I and Mg<sub>3</sub>(PO<sub>4</sub>)<sub>2</sub>-II are monoclinic with space group of *P2<sub>1</sub>/n*, where PO<sub>4</sub><sup>3-</sup> anions in Mg<sub>3</sub>(PO<sub>4</sub>)<sub>2</sub>-I are linked by an identical layer consisting MgO<sub>5</sub> polyhedra and MgO<sub>6</sub> polyhedra, while PO<sub>4</sub><sup>3-</sup> anions in Mg<sub>3</sub>(PO<sub>4</sub>)<sub>2</sub>-II are linked by two inequivalent layers consisting two consistent MgO<sub>6</sub> polyhedra and stacking along the *b*-axis [33–34].

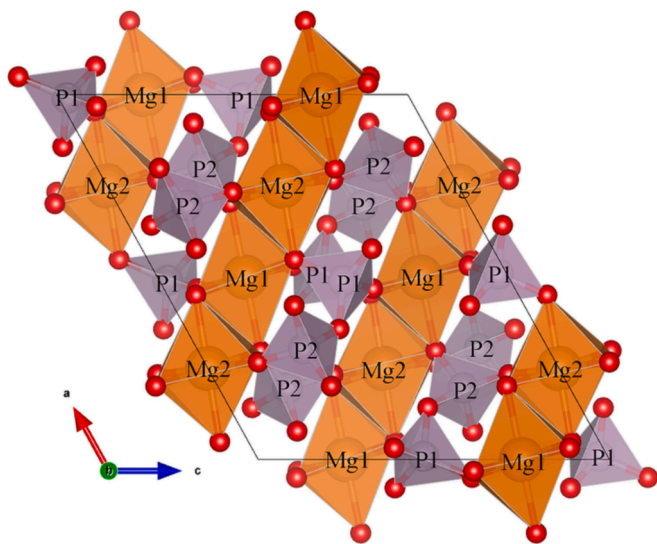


Fig. 1. The crystal structure of MgP<sub>2</sub>O<sub>6</sub>. The red balls represent oxygen. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

**Table 1**  
 Constants determined in  $\nu_i = \nu_0 + b_i T$  of  $\text{MgP}_2\text{O}_6$  at ambient pressure.

No.	$\nu_{i0}$	$-b_i \times 10^2$	$R^2$	$\gamma_{iP}$	$b_i/\nu_{i0} \times 10^5$	$\nu_{297}$	Assignment
1	1355.3(3)	2.61(4)	0.995	0.60	-1.93	1347.0(1)	} $\nu_{as}(\text{PO}_2)$
2	1330.7(1)	2.70(2)	0.999	0.63	-2.03	1322.6(1)	
3	1302.3(3)	2.10(9)	0.988	0.50	-1.61	1296.2(1)	
4	1294.8(3)	2.16(5)	0.992	0.52	-1.67	1288.1(1)	} $\nu_s(\text{PO}_2)$
5	1172.9(2)	1.82(3)	0.994	0.48	-1.55	1166.9(1)	
6	1153.0(1)	1.66(2)	0.997	0.45	-1.44	1147.9(1)	} $\nu_{as}(\text{POP})$
7	1067.8(5)	2.08(7)	0.980	0.61	-1.95	1060.9(2)	
8	938.7(8)	2.51(13)	0.969	0.83	-2.67	930.6(3)	} $\nu_s(\text{POP})$
9	825.8(2)	1.54(3)	0.994	0.58	-1.86	821.2(2)	
10	692.3(2)	1.36(2)	0.996	0.61	-1.96	687.9(1)	} $\text{PO}_2$ bend POP bend
11	620.9(2)	1.54(2)	0.997	0.77	-2.48	616.2(1)	
12	555.7(1)	1.16(2)	0.995	0.65	-2.09	552.2(1)	
13	512.6(2)	1.17(2)	0.993	0.71	-2.28	509.2(4)	
14	463.4(1)	0.75(2)	0.989	0.50	-1.62	461.5(2)	
15	443.0(2)	0.93(4)	0.977	0.65	-2.10	440.6(2)	
16	419.4(1)	0.95(2)	0.994	0.71	-2.27	416.7(1)	
17	400.9(7)	0.49(8)	0.863	0.38	-1.22	399.8(5)	
18	392.2(3)	1.12(8)	0.928	0.89	-2.86	389.2(2)	
19	359.1(1)	0.78(1)	0.997	0.68	-2.17	356.8(1)	
20	339.8(1)	0.57(3)	0.960	0.52	-1.68	338.0(1)	
21	325.8(2)	0.96(2)	0.989	0.92	-2.95	323.5(2)	
22	303.7(2)	2.01(2)	0.998	2.06	-6.62	298.2(3)	
23	251.1(1)	1.43(2)	0.997	1.77	-5.69	247.1(1)	
24	241.9(3)	1.64(6)	0.981	2.11	-6.78	237.2(5)	
25	219.4(3)	1.57(3)	0.993	2.23	-7.16	215.9(3)	
26	198.6(1)	1.69(1)	0.999	2.65	-8.51	193.7(1)	
27	173.5(3)	1.48(3)	0.994	2.66	-8.53	170.8(3)	
28	147.3(1)	1.09(2)	0.997	2.31	-7.40	144.0(2)	
29	100.5(1)	0.55(1)	0.996	1.70	-5.47	99.1(1)	

$\nu_{i0}$  and  $\nu_{297}$  is in  $\text{cm}^{-1}$ , T in K, the constants  $b_i$  and  $b_i/\nu_{i0}$  have the corresponding units.  $R^2$  is the correlation coefficient.

$\nu_{i0}$  and  $\nu_{297}$  is in  $\text{cm}^{-1}$ , T in K, the constants  $b_i$  and  $b_i/\nu_{i0}$  have the corresponding units.  $R^2$  is the correlation coefficient.

$\text{Mg}_3(\text{PO}_4)_2$ -III is triclinic with space group of  $P-1$  consisting of four inequivalent  $\text{PO}_4^{3-}$  anions linked by sharing corners and edges to distorted  $\text{MgO}_4$ ,  $\text{MgO}_5$  and  $\text{MgO}_6$  polyhedra [35].

### 3.1. Temperature-dependent Raman spectra of $\text{MgP}_2\text{O}_6$

The Raman spectra of  $\text{MgP}_2\text{O}_6$  at various temperatures were shown in Fig. 2. Most Raman peaks can be distinguished even at high temperatures. It is clear that the Raman active vibrations shift to lower frequency regions during heating. It is logical since the bond length becomes larger with heating, indicating weaker bond, i.e., small force constant, and consequent lower vibrational frequency according to

Hooke's law. No new Raman peak appears during heating to 1073 K though one or two peaks become undistinguished at high temperatures due to the overlapping or low intensity, indicating no phase transition. As mentioned above, the melting temperature is about 1438 K, which is higher than 1073 K. Therefore, no melting behavior can be observed in this study.

The variations of Raman shift of vibrations at different temperatures for  $\text{MgP}_2\text{O}_6$  were plotted in Fig. 3. The constants derived by linear regressions were listed in Table 1. The temperature coefficients ( $b_i$ ) of vibrations in  $\text{MgP}_2\text{O}_6$  are negative and vary from  $-2.66$  to  $-0.38 \times 10^{-2} \text{ cm}^{-1} \text{ K}^{-1}$ . Generally, the stretching vibrations of  $\text{MgP}_2\text{O}_6$  (including  $\nu_{as}(\text{PO}_2)$ ,  $\nu_s(\text{PO}_2)$ ,  $\nu_{as}(\text{POP})$  and  $\nu_s(\text{POP})$ ) show larger absolute

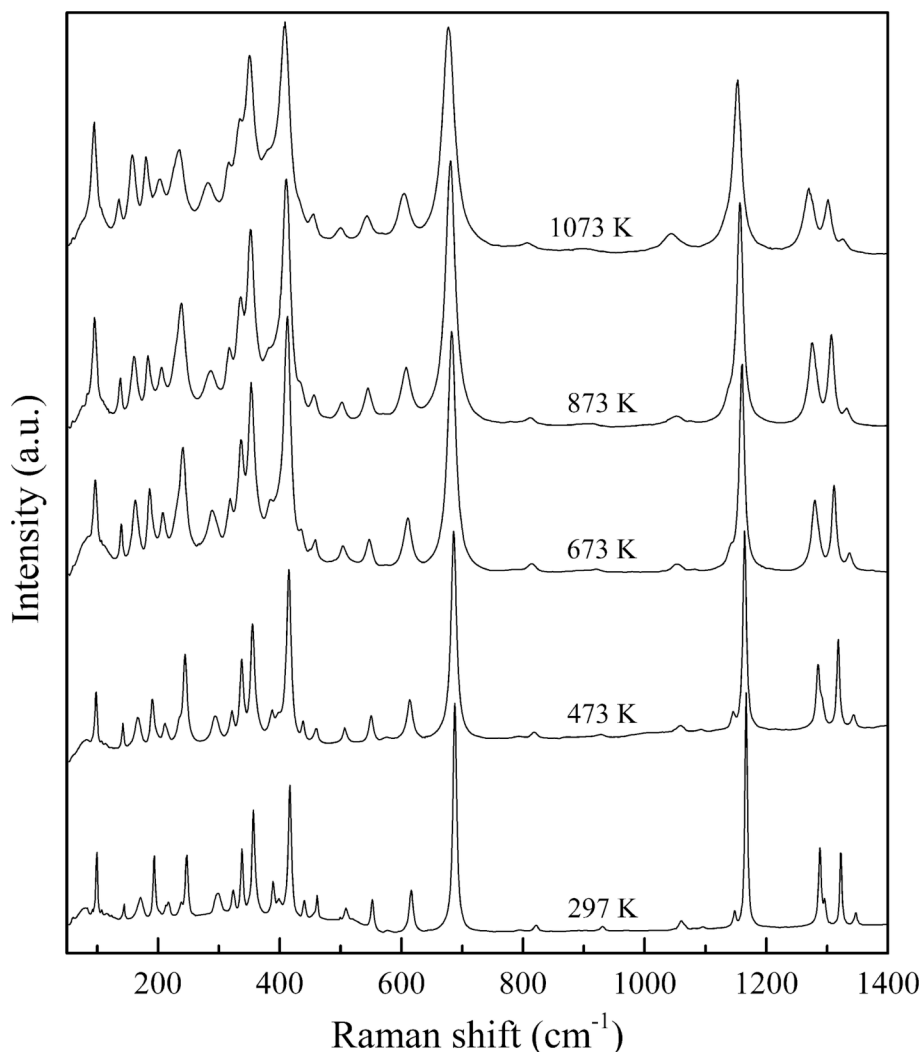


Fig. 2. Typical Raman spectra of  $\text{MgP}_2\text{O}_6$  up to 1073 K at ambient pressure.

temperature-dependence, whereas the  $\text{PO}_2$  and P-O-P bending modes show smaller absolute temperature-dependence. On the other hand, as described in a previous study of Mishra et al. [36], a reduced slope ( $1/\omega_i \partial\omega_i/\partial T$ , as  $b_i/\nu_{i0}$  in Table 1) representing the normalized value of the temperature dependence of mode frequencies can be used to decide the sensitivity of Raman modes with temperature. As listed in Table 1, the calculated reduced slopes are negative and in the range of  $-1.22 \sim -8.53 \times 10^{-5} \text{ K}^{-1}$ . The  $\text{PO}_2$  and POP bending modes show larger reduced slopes, especially the modes less than  $300 \text{ cm}^{-1}$ , than those of stretching vibrations. It means that these bending modes are sensitive to temperature.

### 3.2. Thermal expansion of $\text{MgP}_2\text{O}_6$

Fig. 4 shows the typical XRD patterns of  $\text{MgP}_2\text{O}_6$  from room temperature to 1073 K. With heating, the XRD peaks shift to low  $2\theta$  region. It is reasonable since the  $d$ -spacing expands during heating, resulting in decreasing  $\theta$  according to Bragg's law. Similar to the high-temperature Raman spectroscopic study, there is no new appearing X-ray diffraction peak during heating to 1073 K, implying no phase change. As examples, the refinements of XRD patterns collected at room temperature, 673 K and 1073 K are illustrated in Fig. 5. And the obtained lattice parameters of  $\text{MgP}_2\text{O}_6$  at various temperatures are listed in Table 2. The variations of the unit-cell parameters ( $a$ ,  $b$ ,  $c$ ,  $V$ ) for  $\text{MgP}_2\text{O}_6$  are plotted in Fig. 6.

The thermal expansion coefficient  $\alpha_v = 1/V (\partial V/\partial T)$ , was defined to express the fluctuation of lattice parameters as a function of temperature. If the thermal expansion coefficient  $\alpha_v$  does not change with temperature, integration gives the following expression [37]:

$$\ln(V/V_0) = \alpha_v(T - T_0)$$

In the same way, the axial thermal expansion coefficients along the  $a$ -,  $b$ - and  $c$ -axis can be obtained by the following expressions:

$$\ln(a/a_0) = \alpha_a(T - T_0)$$

$$\ln(b/b_0) = \alpha_b(T - T_0)$$

$$\ln(c/c_0) = \alpha_c(T - T_0).$$

Using the lattice parameters listed in Table 2 and the above expressions, the thermal expansion coefficients of  $\text{MgP}_2\text{O}_6$  are calculated as  $3.21(2) \times 10^{-5} \text{ K}^{-1}$ ,  $1.47(2) \times 10^{-5} \text{ K}^{-1}$ ,  $1.25(5) \times 10^{-5} \text{ K}^{-1}$  and  $0.39(3) \times 10^{-5} \text{ K}^{-1}$  for  $V$ -,  $a$ -,  $b$ - and  $c$ -axis, respectively, as listed in Table 3. Obviously, axial thermal expansion anisotropy exists in  $\text{MgP}_2\text{O}_6$  since the  $a$ - and  $b$ -axis show larger thermal expansion coefficients than that of the  $c$ -axis. It is reasonable since the atomic arrangements along the  $a$ - and  $b$ -axis are looser than those along the  $c$ -axis.

Table 3 also shows the thermal expansion coefficients of other compounds in the  $\text{MgO-P}_2\text{O}_5$  system. Due to different structures and atomic arrangements, these compounds show different bulk and axial

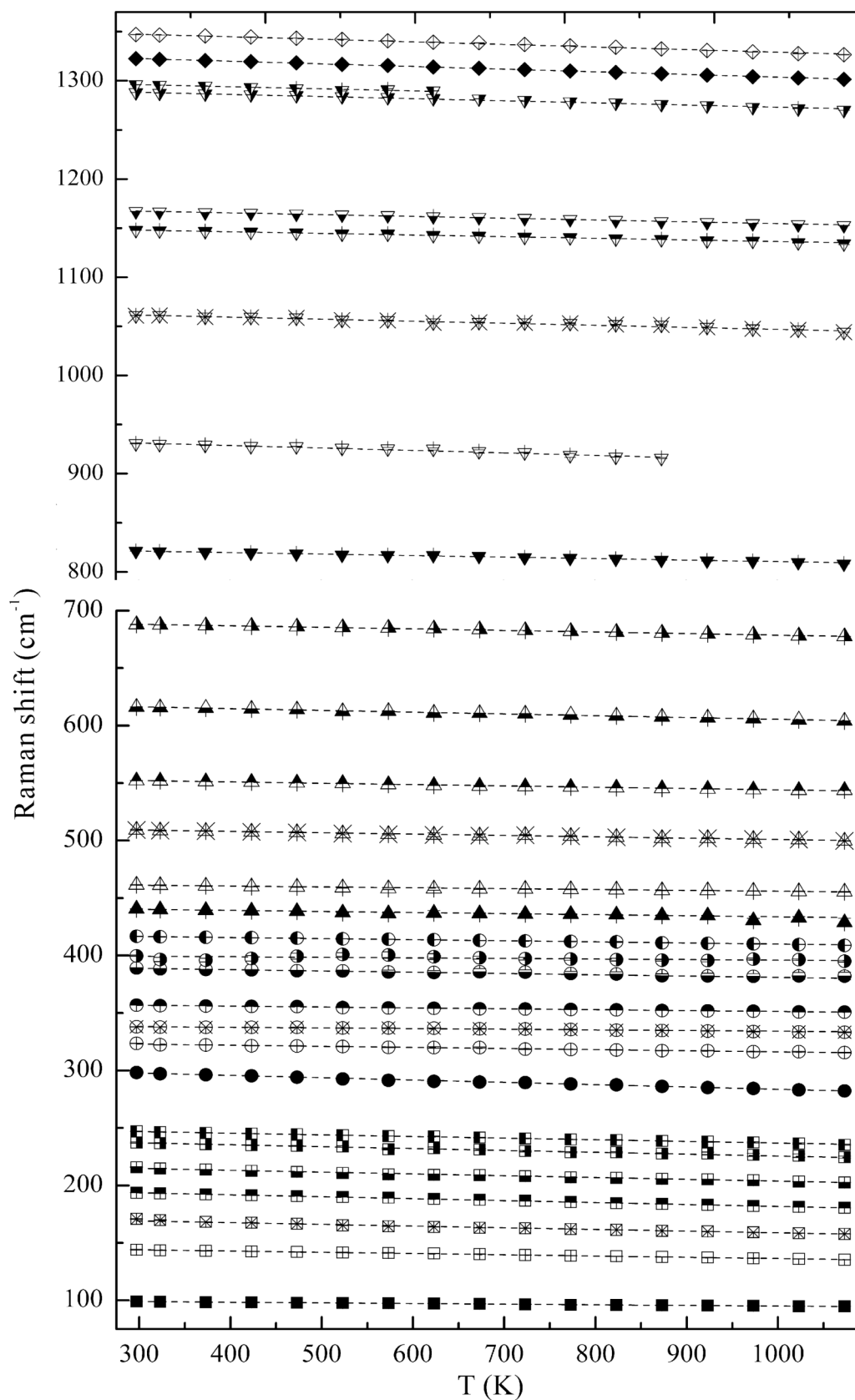


Fig. 3. Temperature dependence of the Raman active vibrations for MgP<sub>2</sub>O<sub>6</sub> at ambient pressure.



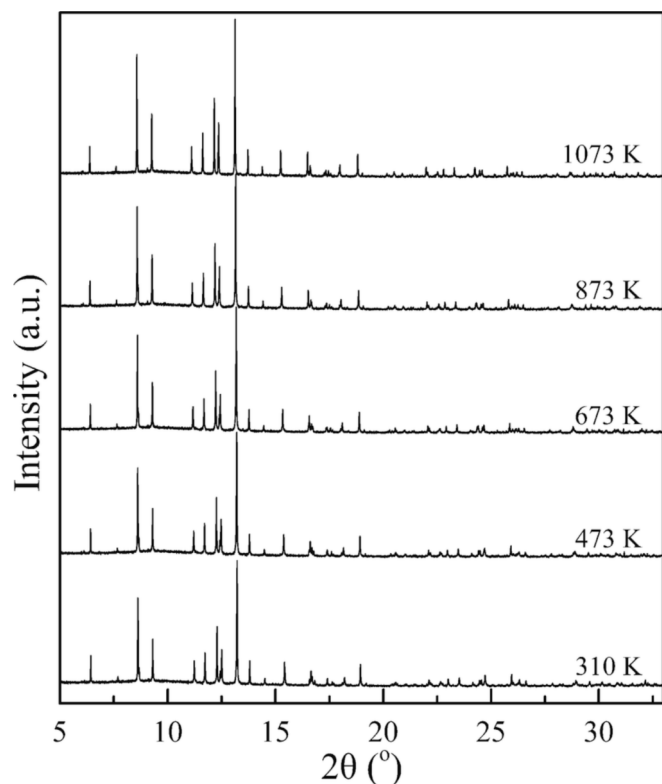


Fig. 4. Typical XRD patterns of  $\text{MgP}_2\text{O}_6$  up to 1073 K at ambient pressure.

thermal expansion behaviors.  $\text{MgP}_2\text{O}_6$  contains  $\text{P}_4\text{O}_{12}^{4-}$  anionic groups and  $\text{MgO}_6$  octahedra,  $\beta\text{-Mg}_2\text{P}_2\text{O}_7$  includes  $\text{P}_2\text{O}_7^{4-}$  anionic groups and  $\text{MgO}_6$  octahedra, and  $\text{Mg}_3(\text{PO}_4)_2$  polymorphs involve  $\text{PO}_4^{3-}$  anionic groups and different Mg-O polyhedra.

### 3.3. Isobaric mode Grüneisen parameters

Based on the temperature-dependent Raman spectra and thermal expansion coefficient, the isobaric mode Grüneisen parameters ( $\gamma_{iP}$ ) of  $\text{MgP}_2\text{O}_6$  can be calculated by using the following formula [38–39]:

$$\gamma_{iP} = -\frac{1}{\alpha} \left( \frac{\partial \ln \nu_i}{\partial T} \right)_P$$

where  $\alpha$  is the volume thermal expansion coefficient, as  $3.21(2) \times 10^{-5} \text{ K}^{-1}$  obtained in this study.

The calculated values of  $\gamma_{iP}$  for different  $\text{P}_4\text{O}_{12}^{4-}$  vibrational modes of  $\text{MgP}_2\text{O}_6$  are also listed in Table 1, which are in the range of 0.38–2.66, and the averaged isobaric mode Grüneisen parameter is 1.06. It is noted that the mode Grüneisen parameter for the low frequency modes ( $100\text{--}300 \text{ cm}^{-1}$ ) are with high values, as listed in Table 1. Indeed,  $\gamma_{iP} = -b_i/\nu_{i0}/\alpha_v$ . Therefore, the larger reduced slopes of the low frequency modes ( $100\text{--}300 \text{ cm}^{-1}$ ) result in higher mode Grüneisen parameters, which also mean that these modes are sensitive to temperature. In our previous study [23], the  $\gamma_{iP}$  for  $\text{P}_2\text{O}_7^{4-}$  vibrational modes of  $\beta\text{-Mg}_2\text{P}_2\text{O}_7$  are in the range of 0.27–1.86 with an average of 0.77. The  $\gamma_{iP}$  for  $\text{PO}_4^{3-}$  vibrational modes of  $\text{Mg}_3(\text{PO}_4)_2$  polymorphs (I, II, III) are 0.35–0.84, 0.31–0.57, and 0.07–1.33 with averages of 0.54, 0.40, and 0.51, respectively [22]. The discrepancy might be contributed to the structural evolution of the  $\text{PO}_4$  tetrahedra during heating in different anionic groups, i.e.,  $\text{P}_4\text{O}_{12}^{4-}$ ,  $\text{P}_2\text{O}_7^{4-}$  and  $\text{PO}_4^{3-}$ , which is also related to the crystal structure (arrangement of atoms) and thermal expansion.

## 4. Conclusions

In this study, the Raman spectra and X-ray diffraction patterns of

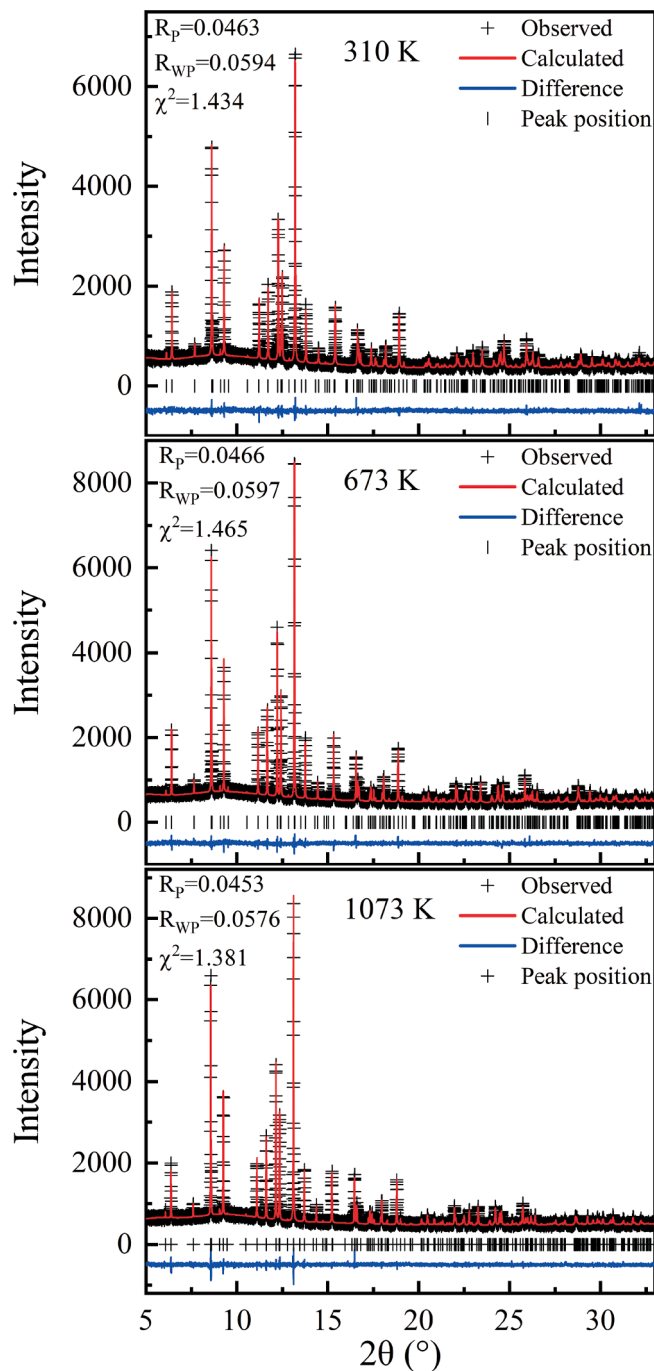


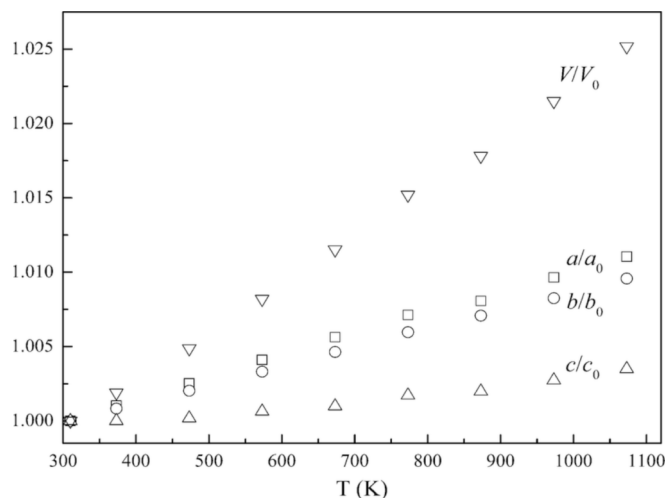
Fig. 5. Refined X-ray diffraction patterns of  $\text{MgP}_2\text{O}_6$  collected at room temperature, 673 K and 1073 K. Black crosses represent the observed XRD pattern, red line, bottom blue curve and red bars are the calculated XRD pattern, the difference and peak positions determined by Rietveld analysis. The respective  $R_p$ ,  $R_{wp}$  and  $\chi^2$  values were also included. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

$\text{MgP}_2\text{O}_6$  were measured from room temperature to 1073 K at ambient pressure. No temperature-induced phase transformation was observed. The Raman shifts of all observed vibrations for  $\text{MgP}_2\text{O}_6$  continuously and linearly decrease with increasing temperature in the dependences of  $-2.61 \times 10^{-2} \sim -0.49 \times 10^{-2} \text{ cm}^{-1} \text{ K}^{-1}$ . The thermal expansion coefficient of  $\text{MgP}_2\text{O}_6$  was determined as  $3.21(2) \times 10^{-5} \text{ K}^{-1}$  and an axial anisotropic thermal expansion exists since the thermal expansion coefficients for  $a$ -,  $b$ - and  $c$ -axis are  $1.47(2) \times 10^{-5} \text{ K}^{-1}$ ,  $1.25(5) \times 10^{-5} \text{ K}^{-1}$

**Table 2**Lattice parameters of MgP<sub>2</sub>O<sub>6</sub> at various temperatures and ambient pressure.

T (K)	a (Å)	b (Å)	c (Å)	β (°)	V (Å <sup>3</sup> )
310	11.7324(2)	8.2628(2)	9.9251(2)	118.905(1)	842.3(5)
373	11.7446(2)	8.2697(2)	9.9252(2)	118.905(1)	843.9(1)
473	11.7620(1)	8.2796(1)	9.9269(1)	118.898(1)	846.4(1)
573	11.7806(1)	8.2902(1)	9.9315(1)	118.889(1)	849.2(1)
673	11.7985(1)	8.3011(1)	9.9350(1)	118.880(1)	852.0(1)
773	11.8160(1)	8.3121(1)	9.9422(1)	118.866(1)	855.1(1)
873	11.8271(2)	8.3213(1)	9.9449(1)	118.844(1)	857.3(1)
973	11.8456(1)	8.3310(1)	9.9523(1)	118.834(1)	860.4(1)
1073	11.8620(1)	8.3419(1)	9.9598(1)	118.819(1)	863.5(1)

Standard deviations are in parentheses.

**Fig. 6.** Relative variations of unit-cell parameters with temperature for MgP<sub>2</sub>O<sub>6</sub>.**Table 3**Thermal expansion ( $\times 10^{-5} \text{ K}^{-1}$ ) of compounds in the MgO-P<sub>2</sub>O<sub>5</sub> system.

Compound	Space group	$a_a$	$a_b$	$a_c$	$a_v$	Ref.
MgP <sub>2</sub> O <sub>6</sub>	C2/c	1.47 (2)	1.25 (5)	0.39 (3)	3.21 (2)	This study
β-MgP <sub>2</sub> O <sub>7</sub>	C2/m	1.48 (7)	1.02 (1)	0.78 (3)	2.97 (8)	Xue et al. (2022)
Mg <sub>3</sub> (PO <sub>4</sub> ) <sub>2</sub> -I	P2 <sub>1</sub> /n	1.04 (1)	2.06 (3)	0.21 (3)	3.31 (4)	Hu et al. (2022)
Mg <sub>3</sub> (PO <sub>4</sub> ) <sub>2</sub> -II	P2 <sub>1</sub> /n	1.42 (1)	1.34 (1)	1.16 (1)	3.91 (4)	Hu et al. (2022)
Mg <sub>3</sub> (PO <sub>4</sub> ) <sub>2</sub> -III	P-1	1.09 (1)	1.57 (2)	1.21 (2)	3.25 (5)	Hu et al. (2022)

Standard deviations are in parentheses.

and  $0.39(3) \times 10^{-5} \text{ K}^{-1}$ , respectively. The isobaric mode Grüneisen parameters of MgP<sub>2</sub>O<sub>6</sub> were calculated in the range of 0.38–2.66 with an average of 1.06.

### CRediT authorship contribution statement

**Weihong Xue:** Investigation, Writing – original draft. **Bo Dai:** Data curation. **Kuan Zhai:** Data curation. **Haipeng Song:** Data curation. **Wen Wen:** Methodology, Resources, Funding acquisition. **Shuangmeng Zhai:** Conceptualization, Writing – review & editing, Funding acquisition.

### Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

### Data availability

Data will be made available on request.

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